

# Mordell-Weil Torsion in the Mirror of Multi-Sections

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Based on

- [arXiv:1408.4808](#) with: D. Klevers, D. Mayorga, H. Piragua and J. Reuter
- [arXiv:1604.00011](#) with: J. Reuter and T. Schimannek

Regional Meeting 2016, Blacksburg  
April 23rd 2016



- 1 F-theory, singular fibers and sections
- 2 Torus Fibers in 2D Ambient Spaces
- 3 Generalizing to 3D Ambient Spaces
- 4 Example: Nef  $(122, 0)$
- 5 Conclusion and Outlook

# The F-theory picture

F-theory: Take the Type IIB axio dilaton:  $\tau = C_0 + ig_s^{-1}$  with

- Theory invariant under  $\tau \rightarrow \frac{a\tau+b}{c\tau+d}$  with  $M = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \text{SL}(2, \mathbb{Z})$
- In addition the  $C_2$  and  $B_2$  must transform as a doublet:

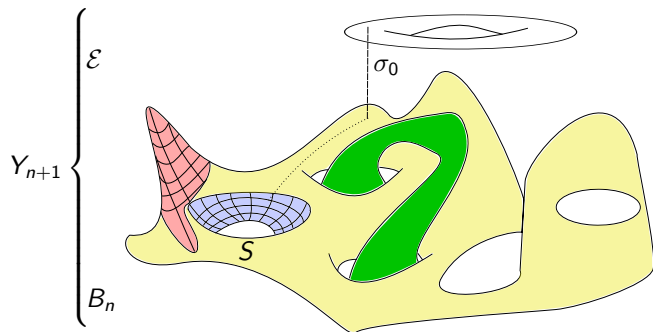
$$\begin{pmatrix} C_2 \\ B_2 \end{pmatrix} \rightarrow M \begin{pmatrix} C_2 \\ B_2 \end{pmatrix} = \begin{pmatrix} aC_2 + B_2 \\ cC_2 + dB_2 \end{pmatrix},$$

- We interpret this structure as coming from the geometry of a torus  $\mathcal{E}$
- The full geometry is a *torus*-fibered n-fold  $Y_n$

$$\begin{array}{ccc} \mathcal{E} & \rightarrow & Y_{n+1} \\ & & \downarrow \\ & & B_n \end{array}$$

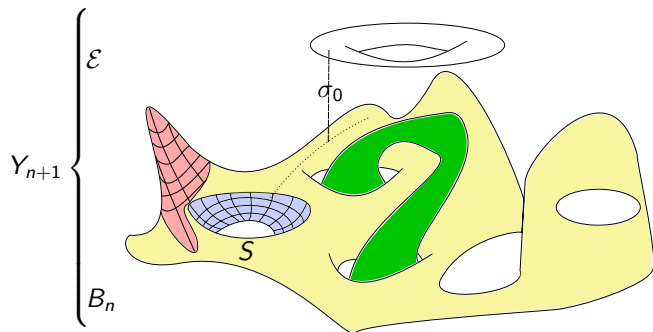
- In the M-theory dual picture, the F-theory fiber Volume is taken to zero, only the Base  $B_n$  is physical

# F-theory Geometry



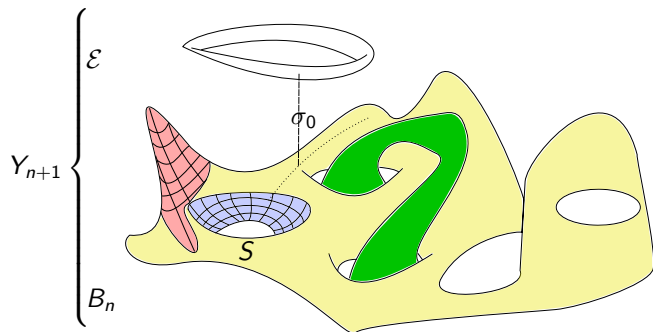
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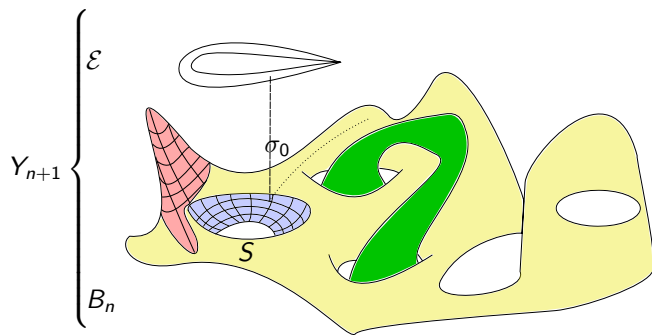
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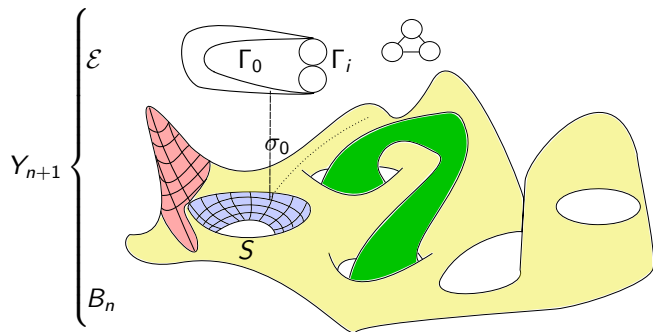


- The zero section  $\sigma_0$  tracks the varying fiber over every point in the base  $B_n$

Over certain codimension 1 (or higher) loci the fiber degenerates:

- $\tau \rightarrow \infty$  : @ D7-brane divisor in the base  $S$
- The whole CY geometry singular resolution in the fiber required

# F-theory Geometry

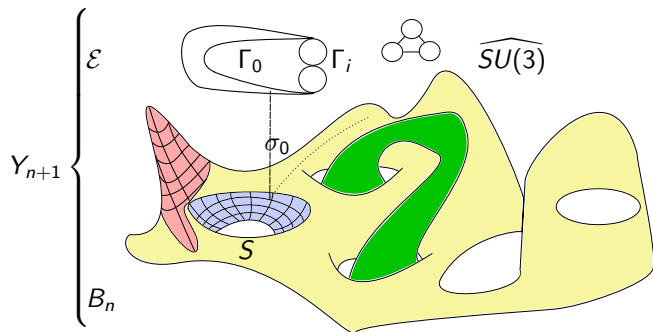


The singularity is of ADE type and can be resolved by gluing in a tree of  $\mathbb{P}^1$ 's

- The  $\mathbb{P}^1$ 's introduce
  - cycles:  $\Gamma_i$  and divisors:  $D_i$  with  $i = 1 \dots n$
  - intersecting  $\Gamma_i \cdot D_j = -C_{ij}$
- The zero point  $\sigma_0$  identifies the affine Node  $\Gamma_0$
- $C_{i,j}$  the Cartan matrix of an affine Lie-group.



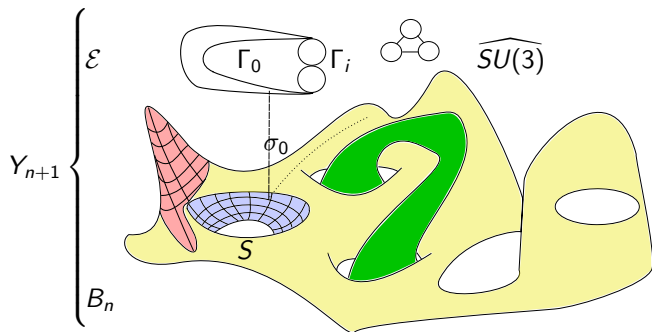
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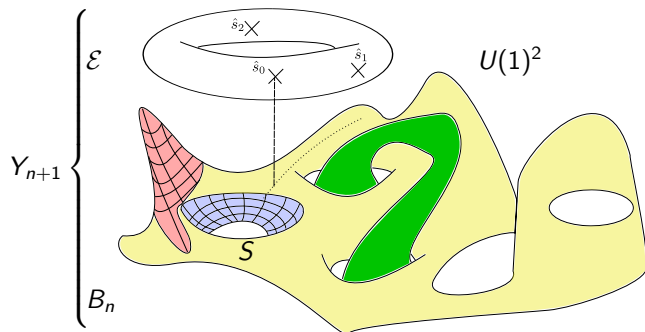
Three sources of Vector-multiplets:

- Cartan-Vector multiplets from the M-theory  $C_3$ -form reduction:

$$\int_{\Gamma_i} C_3 = A_1^i$$

- M2 branes wrap  $\Gamma_i \cup \Gamma_{i+..}$  that become the massless  $W$  bosons when we take the F-theory limit.

# F-theory Geometry

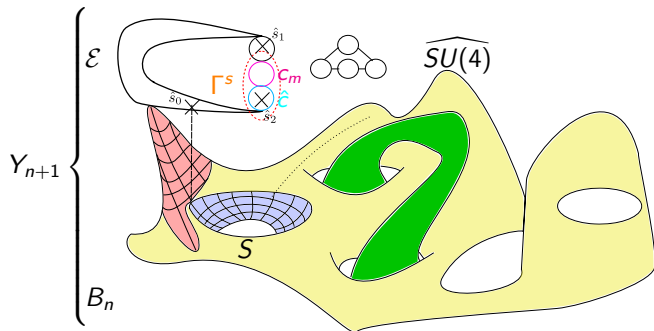


Three sources of Vector-multiplets:

- Additional rational sections  $s_j$  generate the *free Mordell-Weil group* with  $s_j \cdot \mathcal{E} = 1$  give the Shioda map [Vafa Morrison, Morrison Park...]

$$\sigma_i = s_j - s_0 + \pi(K_B^{-1}) + \sum a_n D_n, \quad C_3 = \sigma_j \wedge A_1^j$$

# F-theory Geometry



At codimension two, the fiber degenerates further

- Matter curves from codimension two splits of the curve  $\mathcal{E}$ :

$$\Gamma^s \rightarrow C_m + \hat{C}, \text{ with weights: } (\lambda_i, q_i) = (D_i, \sigma_i) \cdot C_m.$$

Thanks to the splitting of fiber and base there are two general trends to classify F-theory compactifications:

## 1 Classification of fibers $\mathcal{E}$

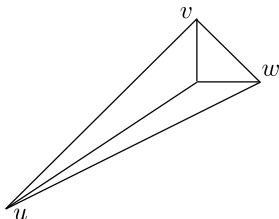
- The Kodaira classification of singular fibers [Kodaira]
- Tate's algorithm [Tate, Katz, Morrison, Schäfer-Nameki, Sully..]
- Classification of fibers in various ambient spaces [Braun, Grimm, Keitel; Klevers, Pena, Piragua, O., Reuter]

## 2 Classification of bases $B_n$

- Classification of two dimensional bases [Morrison, Taylor]
- Classification of non-higgsable clusters [Morrison, Taylor; Grassi, Halverson, Shaneson, Taylor]

## 3 **Bonus ingredient in 4D:** Classification of fluxes [Bizet, Klemm, Lopes; A. Braun Watari ]

# The Fiber description



## Canonical choice: Weierstrass Form

The Weierstrass form as vanishing degree six polynomial  $P_{(1,3,2)}[6]$  in  $[u, v, w]$ :

$$v^2 - w^3 - f(b)wu^4 - g(b)u^6 = 0, \quad \Delta = 27g^2 + 4f^3$$

- **zero section:**  $\sigma_0: [u, v, w] = [0, 1, 1]$
- **Base Dependency** in only two sections  $f(b), f(b)$
- **Discriminant:**  $\Delta = 0 \rightarrow$  singularity directly visible  
 $\rightarrow$  Classification of all codimension 1,2,3 ideals  $V_i$ : with  $\Delta_i = 0$  possible

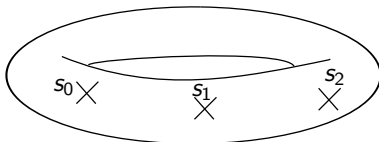
# The Kodaira Classification

ord(f)	ord(g)	ord( $\Delta$ )	Fiber Type	Singularity Type	group
$\geq 0$	$\geq 0$	0	smooth	none	-
0	0	$n$	$I_n$	$A_{n-1}$	$SU(n)$
$\geq 1$	1	2	$II$	none	-
1	$\geq 2$	3	$III$	$A_1$	$SU(2)$
$\geq 2$	2	4	$IV$	$A_2$	$SU(3)$
2	$\geq 3$	$n+6$	$I_n^*$	$D_n+4$	$SO(2n+8)$
$\geq 2$	3	$n+6$	$I_n^*$	$D_n+4$	$SO(2n+8)$
$\geq 3$	4	8	$IV^*$	$E_6$	$E_6$
3	$\geq 5$	9	$III^*$	$E_7$	$E_7$
$\geq 4$	5	10	$II^*$	$E_8$	$E_8$
			$I_1$	none	$U(1)^n ?$
			$I_1$	none	$Z_{n_j} ?$
			$I_1$	none	$1/Z_{n_j} ?$

$$v^2 - w^3 - f(b)wu^4 - f(b)u^6 = 0, \quad \Delta = 27g^2 + 4f^3$$

- What about  $U(1)$ , discrete and quotient symmetries?

# The Zoo of Sections

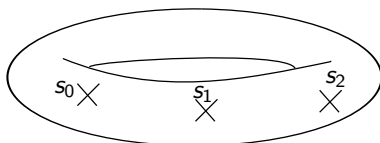


## Rational Sections

Rational sections intersect the fiber once  $s_i \cdot \mathcal{E} = 1$ . They form the *Mordell-Weil group* and geometric addition with the zero-section as neutral element



# The Zoo of Sections



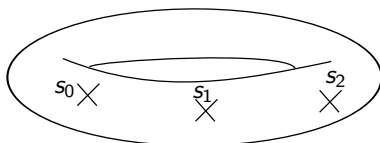
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- 1 **free part**  $\mathbb{Z}^r$ :  $U(1)^r$  symmetries are obtained from the Shioda map, a vertical divisor obtained from the section [Morrison, Park]

$$\sigma_i = s_j - s_0 + \pi(K_B^{-1}) + \sum a_n D_n,$$

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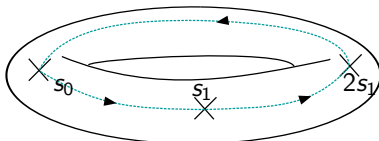
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- 2 **torsion**  $\mathbb{Z}_n$  if sections  $n \cdot s_i \sim s_0$ . Generator of  $\mathbb{Z}_n$  quotient symmetries  $G = G' / \mathbb{Z}_n$ .

Torsion Shioda map  $\sigma_T^{(n)}$  that is a trivial  $\mathbb{Q}$ -divisor  $n \cdot \sigma_T^{(n)} \sim 0$

Matter curves must have charge  $\sigma_T^{(n)} \cdot m = 0 \pmod n$  [Mayrhofer, Morrison, Till, Weigand]

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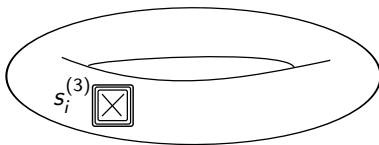
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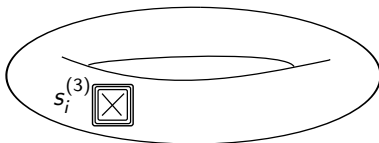
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## Multi-Sections

- A section with  $s_i^{(n)} \cdot \mathcal{E} = n > 1$  is an n-section. If the torus admits only multi-sections it is a genus-one curve.

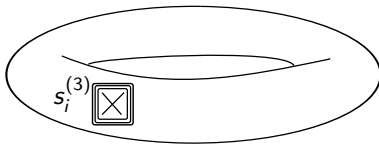
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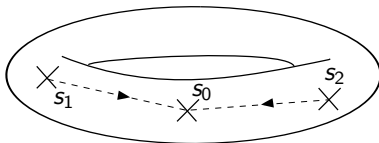
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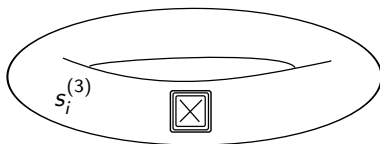
## Multi-Sections

- The  $n$ -section generates a discrete  $\mathbb{Z}_n$  symmetry in the effective field theory
- The  $n$ -section can be obtained by  $n + 1$  collapsing rational sections via a conifold transitions:

In the effective field theory this is a higgsing:  $U(1)^n \rightarrow \mathbb{Z}_n$  [V. Braun, Morrison; Anderson,

Garca-Etxebarria, Grimm, Keitel;]

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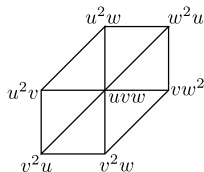
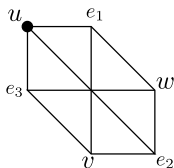
- Why not write down a *discrete Shioda map*: [Klevers, Piragua, Pena, O., Reuter]

$$\sigma_D^{(n)} = s_i^{(n)} + \text{Base} + \sum_i a_i D_i .$$

- Matter charges  $c_m \cdot \sigma_D^{(n)} = k \bmod n$



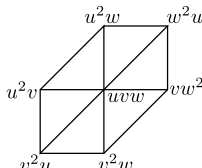
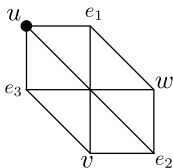
Example:  $U(1)^3 \rightarrow \mathbb{Z}_3$



- Consider a polytope  $\Delta$  and its polar-dual  $\Delta^\circ$

$$P_\Delta = \sum_{m \in \Delta^\circ} \prod_{i=1}^n s_m z_i^{\langle m, \rho_i \rangle + 1} = 0,$$

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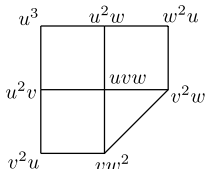
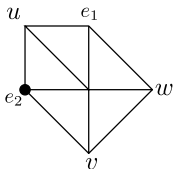
- Start from other ambient spaces such as  $dP_3$ :

$$p_{dP_3} = s_2 e_1 e_3^2 u^2 v + s_3 e_2 e_3^2 u v^2 + s_5 e_1^2 e_3 u^2 w + s_6 e_1 e_2 e_3 u v w + s_7 e_2^2 e_3 v^2 w \\ + s_8 e_1^2 e_2 u w^2 + s_9 e_1 e_2^2 w^2 v.$$

with zero-section  $u \cup p_{dP_3} = 0$  and three sections  $e_i \cup p_{dP_3} = 0$

$U(1)^3$  theory

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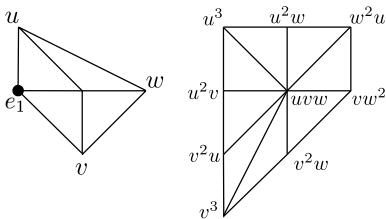
- Collapsing the sections to  $dP_2$

$$p_{dP_2} = s_1 e_2^2 e_1^2 u^3 + s_2 e_2^2 e_1 u^2 v + s_3 e_2^2 u v^2 + s_5 e_2 e_1^2 u^2 w + s_6 e_2 e_1 u v w + s_7 e_2 v^2 w + s_8 e_1^2 u w^2 + s_9 e_1 v w^2,$$

with zero-section  $u \cup p_{dP_2} = 0$  and two sections  $e_i \cup p_{dP_2} = 0$

$U(1)^2$  theory

# Example: $U(1)^3 \rightarrow \mathbb{Z}_3$



- Consider a polytope  $\Delta$  and its polar-dual  $\Delta^\circ$

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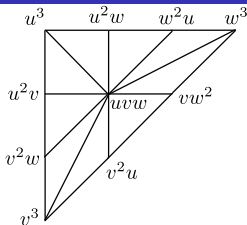
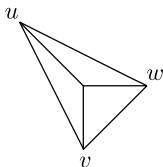
- Collapsing the sections to  $dP_1$

$$p_{P^2} = s_1 u^3 e_1^2 + s_2 u^2 v e_1^2 + s_3 u v^2 e_1^2 + s_4 v^3 e_1^2 + s_5 u^2 w e_1 + s_6 u v w e_1 + s_7 v^2 w e_1 + s_8 u w^2 + s_9 v w^2,$$

with zero-section  $e_1 \cup p_{dP_3} = 0$  and *non-toric* section

$$\boxed{U(1)^1} \text{ theory}$$

# Example: $U(1)^3 \rightarrow \mathbb{Z}_3$



- Consider a polytope  $\Delta$  and its polar-dual  $\Delta^\circ$

$$P_\Delta = \sum_{m \in \Delta^\circ} \prod_{i=1}^n s_m z_i^{\langle m, \rho_i \rangle + 1} = 0,$$

- Collapsing the sections to  $P^2$

$$p_{dP_1} = s_1 u^3 + s_2 u^2 v + s_3 u v^2 + s_4 v^3 + s_5 u^2 w + s_6 u v w \\ + s_7 v^2 w + s_8 u w^2 + s_9 v w^2 + s_{10} w^3,$$

with three-sections only  $(u, v, w) \cup p_{P^2} = 3$

$\boxed{\mathbb{Z}_3}$  theory

# Take Home Message

- Multi-Sections in genus one curves can be thought of as **collapsed rational sections** generating a free Mordell-Weil group
- Rational Sections can be obtained by non ADE blow-ups of the fiber ambient space
- If you look for sections, go search in other fiber ambient spaces

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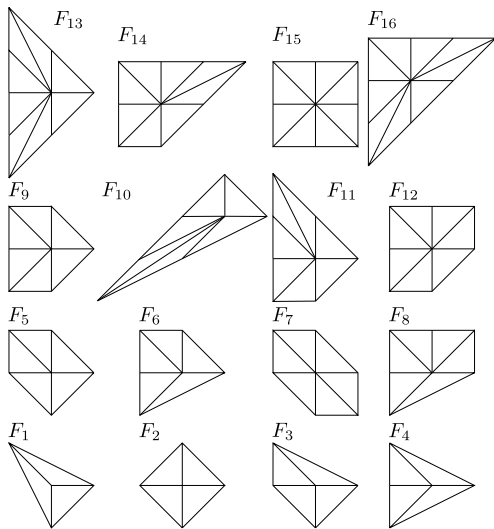
[Kreutzer, Skarke]

- In two dimension there are only 16 of them!
- In three dimensions there are only 3145 of them!

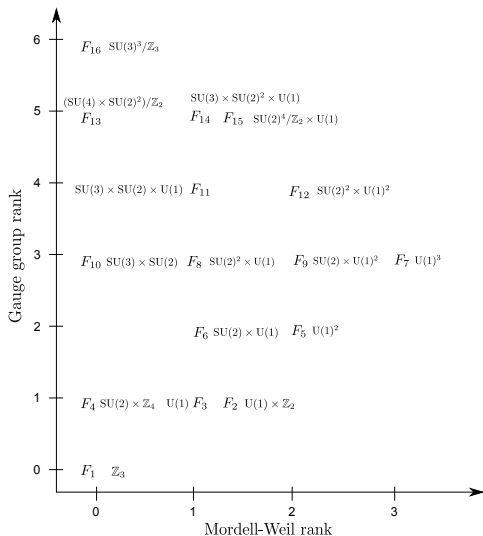
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# The two dimensional Case



# The two dimensional Network



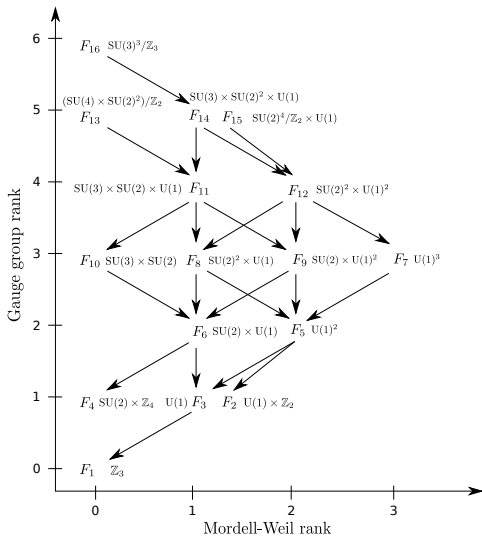
## Torsion

- $\mathbb{Z}_2, \mathbb{Z}_3$  torsion in the upper theories

## Multi Sections

- Genus one curves with two-and three-sections

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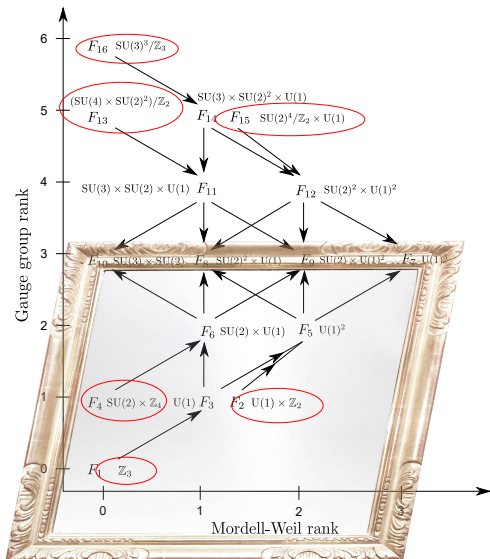
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## Mirror Duality

- $MW(P_\Delta) = MW(P_{\Delta^\circ})$
- $\text{TorMW} \leftrightarrow n\text{-Sections}$

## Multi Sections

- Genus one curves with two-and three-sections

# Understanding the Duality

- Can we relate the appearance of torsion directly from properties of the ambient polytope?
- Can we relate the appearance of (only) multi-sections directly from properties of the ambient polytope?
- Are they connected by  $\Delta \leftrightarrow \Delta^\circ$  ?
- Does this generalize?

## Torsion and Multi-Sections from lattice refinement

- Remember: Lattice points  $\rho_i \in \Delta$  correspond to divisors  $D_i$ . From a dual lattice point  $m$  we can obtain linear equivalence relations between divisors  $D_i$

$$\sum_i \langle \rho_i, m \rangle D_i \sim 0.$$

# Torsion and Multi-Sections from lattice refinement

- Remember: Lattice points  $\rho_i \in \Delta$  correspond to divisors  $D_i$ . From a dual lattice point  $m$  we can obtain linear equivalence relations between divisors  $D_i$

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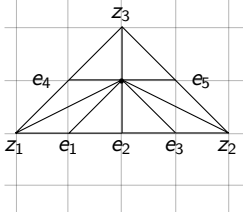
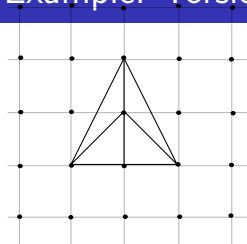
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- The vertices  $\rho_i^{(ver)}$  span a lattice of finite index  $k$  in  $\Delta$
- As  $\Delta$  is of index  $k$ , the dual lattice  $M$  can be refined by a factor of  $k$ .
- The torsion Shioda map is the divisor obtained from the refined dual lattice points

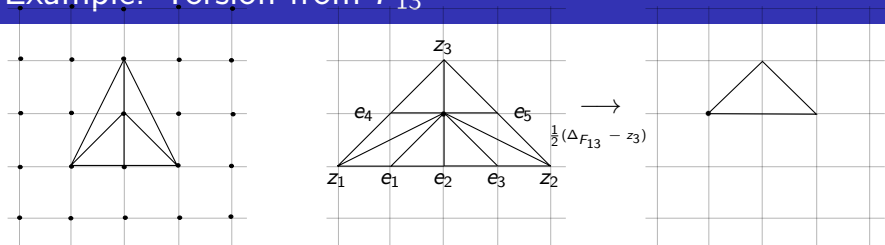
## Example: Torsion from $F_{13}$



### Obtaining the Torsion Shioda map

- The  $F_{13}$  fiber has a generic  $SU(4) \times SU(2)^2/\mathbb{Z}_2$  gauge symmetry that are resolved by the  $e_i$  divisors

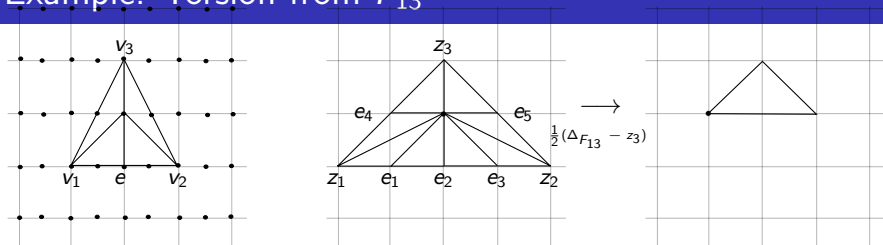
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- There exists a refined dual lattice  $M \rightarrow M'$  with i.e.  $m' = (\frac{1}{2}, 0)$
- Construct the torsion Shioda map from refined lattice point

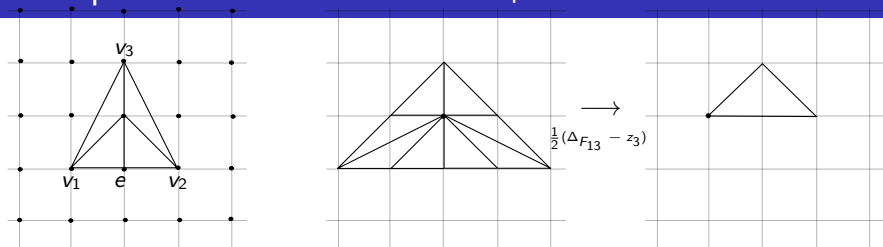
$$\sum_i \langle m', \rho_i \rangle D_i = \sigma_t^{(2)} = [z_2] - [z_1] + \frac{1}{2} (-[e_1] + [e_3] + [e_4] - [e_5]) \quad \checkmark$$

## Discrete Shioda map from lattice refinement

Now we reverse the argument: We have assumed the vertices of  $\rho_i^{(ver)}$  of  $\Delta$  span a lattice of index  $k$  in  $N$

- Now  $\Delta \rightarrow \Delta^\circ$  is a lattice polytope of index  $k$  in  $M$  that corresponds to a divisor i.e. the anticanonical divisor
- The divisor class constructed from polytope  $\Delta^\circ$  is shift invariant  
 $[D_{\Delta^\circ+v}] = [D_{\Delta^\circ}]$
- Scaling a polytope  $\Delta^\circ$  scales the divisor class  $[D_{k\Delta^\circ}] = [kD_{\Delta^\circ}]$
- However the divisor  $D_{\Delta^\circ}$  corresponds a divisor in the anticannonical class of the ambient space.
- Hence the anti-cannonical class is a  $k$ -multiple of an integral class because by assumption is spans a lattice of index  $k$
- Intersections with the class of the vertices in  $\Delta$  are therefore a  $k$ -multiple only

## Example: Multi-Sections from $F_4$



### Obtaining the Torsion Shioda map

We have already seen that we can shrink the dual polytope to

$$\Delta_{F_{13}} \rightarrow \Delta' = \frac{1}{2}(\Delta - w)$$

- Within  $F_4$ ,  $F_{13}$  describes exactly the anticanonical divisor in

$$-K = [v_1] + [v_2] + [v_3] + [e] \sim 2 \underbrace{(2[v_1] + e)}_{D_{\Delta'}}$$

- The fiber in  $F_4$  is a genus one curve only with 2-sections i.e. a theory with  $SU(2) \times \mathbb{Z}_2$  gauge symmetry

# The mirror Conjecture

- This procedure holds true for all 16 fibers constructed from 2D ambient spaces:

**Conjecture 1.** *Given a genus-one fiber  $\mathcal{C}$  for which the Mordell-Weil group of the Jacobian contains torsion, the mirror dual is a genus-one fiber  $\mathcal{C}'$  without a section and vice versa.*



# The mirror Conjecture

Is this phenomenon restricted to fibers in 2D ambient spaces?

## Specialities in 2D

- Vertices have positive intersection with the elliptic curve
  - Divisors are also curves
  - Singularities are torically resolved
  - $\text{Vol}(\Delta) + \text{Vol}(\Delta^\circ) = \text{const}$
  - $\#Points(\Delta) + \#Points(\Delta^\circ) = 12$
- 
- Is there any reason why this should hold in general?
  - In complete intersection fibers, the above constraints do not hold!

- 1 F-theory, singular fibers and sections
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## Recap: Fibers in 3D ambient spaces

- A complete intersection Calabi-Yau is described in terms of a nef partition of a polytope a d-dimensional polytope  $\Delta$

$$\begin{aligned}\Delta &= \Delta_1 + \dots + \Delta_n, & \Delta^\circ &= \langle \nabla_1, \dots, \nabla_n \rangle, \\ \nabla^\circ &= \langle \Delta_1, \dots, \Delta_n \rangle, & \nabla &= \nabla_1 + \dots + \nabla_n,\end{aligned}\tag{1}$$

- This specifies a codimension n Calabi-Yau in a d-dimensional polytope via the intersections of

$$P_{\Delta_i} = \sum_{m \in \Delta_i} \prod_{i=1}^n s_m z_i^{\langle m, \rho_i \rangle + a_i} \in \mathbb{P}_{\Delta^\circ}$$

with  $a_i \geq -1$

- The Mirror CY is cut out by  $P_{\nabla_i} \in \mathbb{P}_{\nabla^\circ}$
- Note: One ambient space can have multiple nef partitions whose mirror dual do not have to live in the same ambient spaces!

# How to check it?

The arguments in 2D do not readily apply. The existence of torsion does not imply multisections in the dual geometry anymore!

In [Braun, Grimm, Keitel '15] 3145 codimension two nef partitions in 3D ambient spaces have been considered

- They constructed *the toric MW group*
- They mapped all curves into their Jacobian Form i.e. into WSF form

A new hope:

- The fiber in  $\mathbb{P}^3$  only admits four sections
- The dual Nef partition with Palp id (3145,0) has  $\mathbb{Z}_4$  Mordell-Weil torsion

Conjecture counter: 2/3145

# Does it generalize?

- We went through the full list again and obtained all intersections of *toric divisors* with the elliptic curve
- Combining this information with the Mordell-Weil group, we indeed get a match in

Conjecture counter: 3086/3145

- So what is wrong about the rest?

# Reducibility and non-toric sections

## Nef partition (4, 0)

Specified by the nef partition

$$\nabla_1 = \langle z_0, z_3, z_4, z_6 \rangle, \quad \nabla_2 = \langle z_1, z_2, z_5, z_8 \rangle,$$

in the  $\mathbb{P}^1 \times \mathbb{P}^2$  ambient space:

$z_0$	$z_1$	$z_2$	$z_3$	$z_4$
$\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$	$\begin{pmatrix} -1 \\ 0 \\ 0 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$	$\begin{pmatrix} 0 \\ -1 \\ -1 \end{pmatrix}$

- Two inequivalent divisor classes:  $[z_0] \sim [z_1]$ ,  $[z_2] \sim [z_3] \sim [z_4]$
- Intersections  $\mathcal{E} \cdot [z_0] = 2$  and  $\mathcal{E} \cdot [z_2] = 3$ .
- The dual nef partition (3013, 1) has no torsion.
- This is not a genus one curve: Construct a non-toric zero-section [Braun, Grimm, Keitel

14]

$$[s_0] = [z_2] - [z_0] + H_B$$

# Reducibility and multi-mirrors

- Lets consider the elliptic curve with CICY equation of (3013, 1) in  $z_i$

$$\begin{aligned} p_1 &= a_0 z_5 z_6 z_7 z_8 z_9 z_{10} z_{11} z_{12} z_{13} z_{15} z_{16} + a_2 z_0 z_3 z_5^2 z_7^2 z_{11} z_{12} z_{13} \\ &\quad + a_1 z_1 z_4 z_6^2 z_8^2 z_{13} z_{15} z_{16} + a_3 z_0 z_1 z_2 z_5 z_6 z_9 z_{12} z_{16} \\ p_2 &= a_5 z_2 z_9^2 z_{10}^2 z_{11} z_{12} z_{15} z_{16} + a_4 z_3 z_4 z_7 z_8 z_{10} z_{11} z_{13} z_{15} + a_6 z_0 z_1 z_2 z_3 z_4 \end{aligned}$$

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- Lets also consider the elliptic curve with CICY equation of (3013,0) in  $z_i$

$$\begin{aligned} p_1 &= a_0 z_5 z_6 z_7 z_8 z_9 z_{10} z_{11} z_{12} z_{13} z_{15} z_{16} + a_2 z_0 z_3 z_5^2 z_7^2 z_{11} z_{12} z_{13} \\ &\quad + a_1 z_1 z_4 z_6^2 z_8^2 z_{13} z_{15} z_{16} + a_3 z_0 z_1 z_2 z_5 z_6 z_9 z_{12} z_{16} \\ p_2 &= a_5 z_2 z_9^2 z_{10}^2 z_{11} z_{12} z_{15} z_{16} + a_4 z_3 z_4 z_7 z_8 z_{10} z_{11} z_{13} z_{15} + a_6 z_0 z_1 z_2 z_3 z_4 \end{aligned}$$



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- As they come from the same ambient space, they also have exactly the same Stanley-Reisner ideal: These are equivalent elliptic curves!

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- As they come from the same ambient space, they also have exactly the same Stanley-Reisner ideal: These are equivalent elliptic curves!
- Both curves have a zero-section and no torsion
- However, the first has (4, 0) as a mirror dual, the second one (5, 1). Are those equivalent too?

# Reducibility

- Nef (5, 1) admits a toric section with CICY equation in  $z_i$

$$p_1 = a_8 z_0 z_2^2 + a_7 z_0 z_2 z_3 + a_5 z_0 z_3^2 + a_6 z_0 z_2 z_4 + a_4 z_0 z_3 z_4 \\ + a_3 z_0 z_4^2 + a_2 z_1 z_2 + a_1 z_1 z_3 + a_0 z_1 z_4$$

$$p_2 = a_{17} z_0 z_2^2 + a_{16} z_0 z_2 z_3 + a_{14} z_0 z_3^2 + a_{15} z_0 z_2 z_4 \\ + a_{13} z_0 z_3 z_4 + a_{12} z_0 z_4^2 + a_{11} z_1 z_2 + a_{10} z_1 z_3 + a_9 z_1 z_4$$

- The CICY equation of (4,0) is very different and has no toric section:

$$p_1 = a_{11} z_0 z_2^2 + a_5 z_1 z_2^2 + a_{10} z_0 z_2 z_3 + a_4 z_1 z_2 z_3 + a_8 z_0 z_3^2 + a_2 z_1 z_3^2 \\ + a_9 z_0 z_2 z_4 + a_3 z_1 z_2 z_4 + a_7 z_0 z_3 z_4 + a_1 z_1 z_3 z_4 + a_6 z_0 z_4^2 + a_0 z_1 z_4^2$$

$$p_2 = a_{17} z_0 z_2 + a_{14} z_1 z_2 + a_{16} z_0 z_3 + a_{13} z_1 z_3 + a_{15} z_0 z_4 + a_{12} z_1 z_4$$

- Those models share the same Weierstrass models with:

$$(5, 1) : f(a_i), g(a_i) = (4, 0) : f(a_i), g(a_i)$$

After relabeling the sections  $a_i$ !

- Hence the curve in (4, 0) is **equivalent** to (5,1)

Conjecture counter: 3088/3145

# Reducibility of CICY fibers

- By this procedure we find that many fibers have an **equivalent (singular) Weierstrass** description related by a simple **relabeling of the sections**
- The 3145 different nef partitions get reduced to 1024 inequivalent fibers
- We find examples where ADE singularities have toric and non-toric resolutions when realized in different ambient spaces

Conjecture counter: 998/1024

## 26 cases to go

Conjecture counter: 998/1024

- In all cases, the curve and its dual are genus-one curves that only admit multi-sections. → No toric Mordell-Weil group. How to test for the torsion?
- We consider the Weierstrass models of those curves (Jacobian) and show that they are birational equivalent to the Weierstrass form with  $k$ -torsion

[Aspinwall, Morrison '98]

# Mirror genus-one curves

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- We consider the Weierstrass models of those curves (Jacobian) and show that they are birational equivalent to the Weierstrass form with  $k$ -torsion  
[Aspinwall, Morrison '98]
- Those models have torsion sections only in their Jacobian
- In 24 cases the genus one curve has only two-sections and two-torsion, just as their mirrors  $\checkmark$

$\leftarrow$ mirror dual $\rightarrow$	$\leftarrow$ mirror dual $\rightarrow$
$(152, 0)$   $(195, 4)$	$(8, 0)$   $(609, 0)$
$(29, 2)$   $(577, 0)$	$(34, 0)$   $(321, 1)$
$(39, 0)$   $(335, 0)$	$(56, 2)$   $(356, 2)$
$(78, 2)$   $(266, 0)$	$(108, 0)$   $(161, 1)$
$(129, 0)$   $(129, 1)$	$(150, 1)$   $(208, 1)$
$(152, 1)$ self-mirror	$(122, 0)$ self-mirror

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[Aspinwall, Morrison '98]

- in four cases, the degree of the Multi-sections and torsion does not match!

$\leftarrow$  mirror dual  $\rightarrow$

nef (5, 3)	nef (2069, 0)
four-sections no torsion	$\mathbb{Z}_2$ torsion one-sections
nef (21, 1)	nef (488, 0)
four-sections $\mathbb{Z}_2$ torsion	$\mathbb{Z}_2$ torsion two-sections

- These theories have non-toric resolution divisors and are connected via a higgsing.
- Apart from the matching of the degree, the conjecture still holds



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Conjecture counter:  $1024(4)/1024$  ✓

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[Aspinwall, Morrison '98]

**Conjecture 2.** *A curve  $C$  constructed as a complete intersection in a toric ambient space such that all the codimension one loci are torically resolved does exhibit Mordell-Weil torsion of degree  $k$  in the Jacobian if and only if the one dimensional generators of the fan  $\{\rho_i : D_i \cdot C \neq 0\}$  span a sublattice  $\tilde{M} \supset M$  of index  $k$ . Up to base divisors a point  $m \in \tilde{M} \setminus M$  corresponds to a torsion Shioda map  $\sigma_t^{(k)}$  via*

$$\sigma_t^{(k)} = \sum_{\rho_i \in \Sigma(1)} \langle m, \rho_i \rangle \cdot D_i.$$

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# A self-dual genus one-curve

- The nef partition is given by:

$$\nabla_1 = \langle z_0, z_3, z_4, z_6 \rangle, \quad \nabla_2 = \langle z_1, z_2, z_5, z_8 \rangle,$$

- in the ambient space

$z_0$	$z_1$	$z_2$	$z_3$
$\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$	$\begin{pmatrix} -1 \\ -1 \\ -1 \end{pmatrix}$
$z_4$	$z_5$	$z_6$	$z_8$
$\begin{pmatrix} 1 \\ -1 \\ -1 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 2 \\ 1 \end{pmatrix}$	$\begin{pmatrix} 0 \\ -1 \\ -1 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}$

- The CICY equations are given as

$$p_1 = a_2 z_0^2 z_4^2 z_6 + a_4 z_0 z_1 z_5^2 z_8 + a_1 z_0 z_3 z_4 z_6 + a_3 z_0 z_2 z_5 z_8 + a_0 z_3^2 z_6,$$
$$p_2 = a_9 z_0 z_1 z_4^2 z_6 + a_8 z_1^2 z_5^2 z_8 + a_6 z_1 z_3 z_4 z_6 + a_7 z_1 z_2 z_5 z_8 + a_5 z_2^2 z_8.$$

- Intersecting the curve:  $[z_0, z_1, z_2, z_3, z_4, z_5, z_6, z_8] \cdot \mathcal{E} = [0, 0, 2, 2, 2, 2, 0, 0]$
- We have only-two sections: this is a genus-one curve
- The vertices span a lattice of index 2 in  $\mathbb{Z}^3$

# Gauge symmetry

$$\begin{array}{cccc} z_0 & z_1 & z_2 & z_3 \\ \hline \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} & \begin{pmatrix} -1 \\ -1 \\ -1 \end{pmatrix} \\ \\ z_4 & z_5 & z_6 & z_8 \\ \hline \begin{pmatrix} 1 \\ -1 \\ -1 \end{pmatrix} & \begin{pmatrix} 0 \\ 2 \\ 1 \end{pmatrix} & \begin{pmatrix} 0 \\ -1 \\ -1 \end{pmatrix} & \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix} \end{array}$$

- The dual lattice is generated by a finer lattice i.e.

$$m' = \left(\frac{1}{2}, \frac{1}{2}, 0\right), m_1 = (0, 1, 0), m_2 = (0, 0, 1)$$

- From  $m'$  we construct the torsion Shioda map:

$$\sigma_T^{(2)} = [z_5] - [z_3] + \frac{1}{2} ([z_0] + [z_1] - [z_6] + [z_8]) .$$

- Indeed:  $z_0, z_1, z_6, z_8$  are resolution divisors
- The intersection with matter  $\sigma_T^{(2)} \cdot c_m \sim 0$  implies for the weights  $\lambda^i$

$$\lambda_m^1 + \lambda_m^2 - \lambda_m^3 + \lambda_m^4 = 0 \pmod{2}$$

- Only bifundamental matter possible

# Consistency of torsion

We can check for the existence of  $\mathbb{Z}_2$  torsion in the Weierstrass form:

- A Weierstrass model with a  $\mathbb{Z}_2$  torsion points admits the following form of Weierstrass coefficients: [Aspinwall, Morrison' 98]

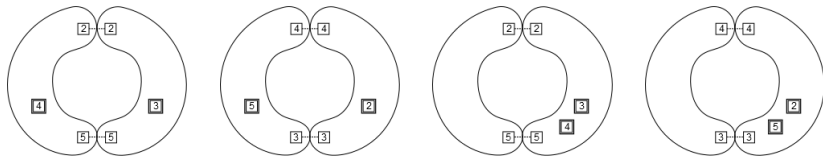
$$f = A_4 - \frac{1}{3}A_2^2, \quad g = \frac{1}{27}A_2(2A_2^2 - 9A_4),$$
$$\Delta = A_4^2(4A_4 - A_2^2),$$

- The  $(122, 0)$  Weierstrass coefficients are related by the birational map:

$$A_2 \rightarrow 4a_1a_4a_5a_6 + a_3^2a_6^2 - 2a_1a_3a_6a_7 + a_1^2a_7^2$$
$$- 4a_0a_2a_7^2 - 4a_1^2a_5a_8 + 16a_0a_2a_5a_8$$
$$- 8a_0a_4a_5a_9 + 4a_0a_3a_7a_9,$$
$$A_4 \rightarrow 16a_0a_5(a_4^2a_5 - a_3a_4a_7 + a_3^2a_8)$$
$$\cdot (a_2a_6^2 - a_1a_6a_9 + a_0a_9^2).$$

- Therefore we have  $\mathbb{Z}_2$  torsion ✓
- We have four  $SU(2)$  singularities in codimension two ✓

# Discrete Shioda map and gauge locus



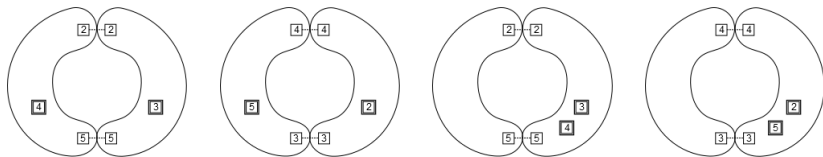
- The CICY equations are

$$p_1 = a_2 z_0^2 z_4^2 z_6 + a_4 z_0 z_1 z_5^2 z_8 + a_1 z_0 z_3 z_4 z_6 + a_3 z_0 z_2 z_5 z_8 + a_0 z_3^2 z_6 ,$$

$$p_2 = a_9 z_0 z_1 z_4^2 z_6 + a_8 z_1^2 z_5^2 z_8 + a_6 z_1 z_3 z_4 z_6 + a_7 z_1 z_2 z_5 z_8 + a_5 z_2^2 z_8 .$$

- We have the following four codimension one,  $I_2$  singularities:

# Discrete Shioda map and gauge locus



- The CICY equations are

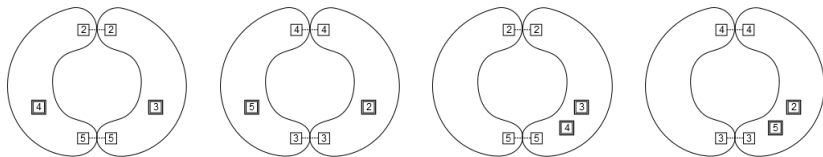
$$p_1 = a_2 z_0^2 z_4^2 z_6 + a_4 z_0 z_1 z_5^2 z_8 + a_1 z_0 z_3 z_4 z_6 + a_3 z_0 z_2 z_5 z_8 + a_0 z_3^2 z_6 ,$$
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- We have the following four codimension one,  $l_2$  singularities:
- toric locus  $l_2$  locus  $L_1 : a_0 = 0$

$$p_1^{(1)} = p' \cdot z_0 ,$$



# Discrete Shioda map and gauge locus



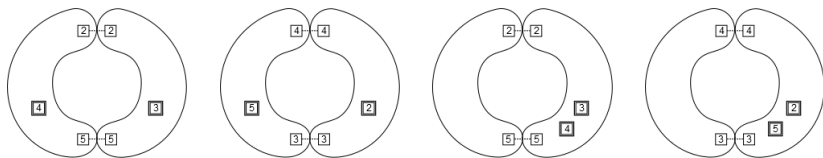
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- We have the following four codimension one,  $l_2$  singularities:
- toric locus  $l_2$  locus  $L_2 : a_5 = 0$

$$p_2 = p_2' \cdot z_1 .$$

# Discrete Shioda map and gauge locus



- The CICY equations are

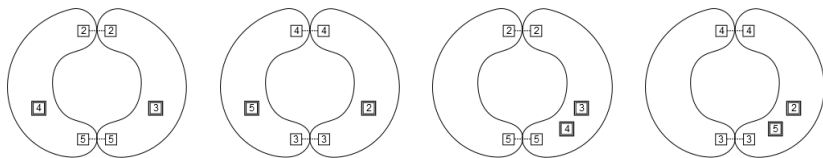
$$p_1 = a_2 z_0^2 z_4^2 z_6 + a_4 z_0 z_1 z_5^2 z_8 + a_1 z_0 z_3 z_4 z_6 + a_3 z_0 z_2 z_5 z_8 + a_0 z_3^2 z_6,$$

$$p_2 = a_9 z_0 z_1 z_4^2 z_6 + a_8 z_1^2 z_5^2 z_8 + a_6 z_1 z_3 z_4 z_6 + a_7 z_1 z_2 z_5 z_8 + a_5 z_2^2 z_8.$$

- We have the following four codimension one,  $I_2$  singularities:
- non-toric  $I_2$  locus  $L_3 : a_4^2 a_5 - a_3 a_4 a_7 + a_3^2 a_8 = 0$  the following combination factorizes:

$$a_3 a_4 z_0 z_5 p_2 - (a_4 a_5 z_2 + a_3 a_8 z_1 z_5) p_1 = p^{(2)} \cdot z_6$$

# Discrete Shioda map and gauge locus



- The CICY equations are

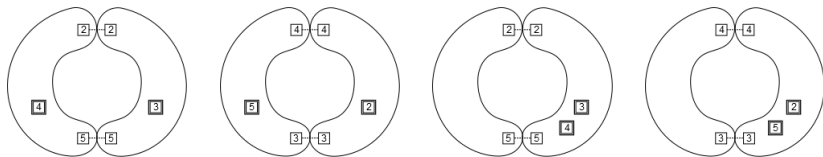
$$p_1 = a_2 z_0^2 z_4^2 z_6 + a_4 z_0 z_1 z_5^2 z_8 + a_1 z_0 z_3 z_4 z_6 + a_3 z_0 z_2 z_5 z_8 + a_0 z_3^2 z_6 ,$$

$$p_2 = a_9 z_0 z_1 z_4^2 z_6 + a_8 z_1^2 z_5^2 z_8 + a_6 z_1 z_3 z_4 z_6 + a_7 z_1 z_2 z_5 z_8 + a_5 z_2^2 z_8 .$$

- We have the following four codimension one,  $I_2$  singularities:
- non-toric  $I_2$  locus  $L_4 : a_2 a_6^2 - a_1 a_6 a_9 + a_0 a_9^2 = 0$  the following combination factorizes:

$$a_6 a_9 z_1 z_4 p_1 - (a_0 a_9 z_3 + a_2 a_6 z_0 z_4) p_2 = p^{(3)} \cdot z_8 ,$$

# Discrete Shioda map and gauge locus



- The CICY equations are

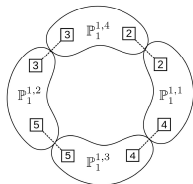
$$p_1 = a_2 z_0^2 z_4^2 z_6 + a_4 z_0 z_1 z_5^2 z_8 + a_1 z_0 z_3 z_4 z_6 + a_3 z_0 z_2 z_5 z_8 + a_0 z_3^2 z_6 ,$$

$$p_2 = a_9 z_0 z_1 z_4^2 z_6 + a_8 z_1^2 z_5^2 z_8 + a_6 z_1 z_3 z_4 z_6 + a_7 z_1 z_2 z_5 z_8 + a_5 z_2^2 z_8 .$$

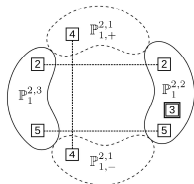
- We have the following four codimension one,  $I_2$  singularities:
- The multisections intersect the fiber and hence we have to orthogonalize the discrete Shioda map;

$$\sigma_{D,4}^{(2)} = [z_4] + \frac{1}{2} ([z_1] + [z_6])$$

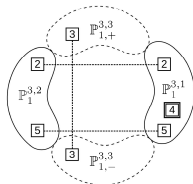
# The matter spectrum



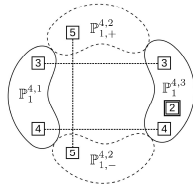
Matter locus 1,  
 $a_0 = a_5 = 0$



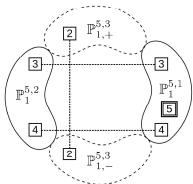
Matter locus 2,  
 $a_0 = a_6 = 0$



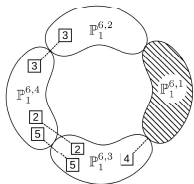
Matter locus 3,  
 $a_0 = a_2 a_6 - a_1 a_9 = 0$



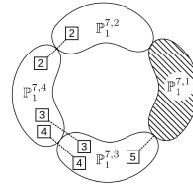
Matter locus 4,  
 $a_5 = a_3 = 0$



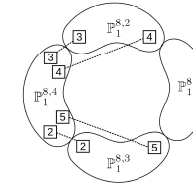
Matter locus 5,  
 $a_5 = a_4 a_7 - a_3 a_8 = 0$



Matter locus 6,  
 $a_0 = 0, L_2$



Matter locus 7,  
 $a_5 = 0, L_3$



Matter locus 8,  
 $L_2, L_3$

# The matter spectrum

Locus	$(f, g, \Delta)$	$SU(2)^4 \times \mathbb{Z}_2$ Rep.
$a_0 = 0, a_5 = 0$	$(0, 0, 4)$	$(2, 2, 1, 1)_{\frac{1}{2}}$
$a_0 = 0, a_6 = 0$	$(0, 0, 4)$	$(2, 1, 1, 2)_1$
$a_0 = 0,$ $a_2 a_6 - a_1 a_9 = 0$	$(0, 0, 4)$	$(2, 1, 1, 2)'_0$
$a_5 = 0, a_3 = 0$	$(0, 0, 4)$	$(1, 2, 2, 1)_0$
$a_5 = 0,$ $a_4 a_7 - a_3 a_8 = 0$	$(0, 0, 4)$	$(1, 2, 2, 1)'_1$
$a_0 = 0,$ $a_4^2 a_5 - a_3 a_4 a_7 + a_3^2 a_8 = 0$	$(0, 0, 4)$	$(2, 1, 2, 1)_{-\frac{1}{2}}$
$a_5 = 0,$ $a_2 a_6^2 - a_1 a_6 a_9 + a_0 a_9^2 = 0$	$(0, 0, 4)$	$(1, 2, 1, 2)_{\frac{1}{2}}$
$a_4^2 a_5 - a_3 a_4 a_7 + a_3^2 a_8 = 0,$ $a_2 a_6^2 - a_1 a_6 a_9 + a_0 a_9^2 = 0$	$(0, 0, 4)$	$(1, 1, 2, 2)_{\frac{1}{2}}$

# The matter spectrum

- Only bifundamental matter ✓
- All matter curves distinguished by a unique quantum number ✓

$$\begin{aligned} &(\mathbf{2}, \mathbf{1}, \mathbf{1}, \mathbf{2})_1, & (\mathbf{2}, \mathbf{1}, \mathbf{1}, \mathbf{2})'_0 \\ &(\mathbf{1}, \mathbf{2}, \mathbf{2}, \mathbf{1})_1, & (\mathbf{1}, \mathbf{2}, \mathbf{2}, \mathbf{1})'_0 \end{aligned}$$

- The  $\mathbb{Z}_2$  also restricts the Yukawa couplings as expected

$$\begin{aligned} Y_1 &: (\mathbf{2}, \mathbf{2}, \mathbf{1}, \mathbf{1})_{\frac{1}{2}} \cdot (\mathbf{2}, \mathbf{1}, \mathbf{2}, \mathbf{1})_{-\frac{1}{2}} \cdot (\mathbf{1}, \mathbf{2}, \mathbf{2}, \mathbf{1})_0, \\ Y_2 &: (\mathbf{2}, \mathbf{2}, \mathbf{1}, \mathbf{1})_{\frac{1}{2}} \cdot \overline{(\mathbf{2}, \mathbf{1}, \mathbf{2}, \mathbf{1})_{-\frac{1}{2}}} \cdot (\mathbf{1}, \mathbf{2}, \mathbf{2}, \mathbf{1})'_1. \end{aligned}$$

- Without the  $\mathbb{Z}_2$  symmetry the geometric different Yukawa couplings would be the same

# Summary and Outlook

- We have given strong evidence, that genus-one curves with multi-sections are mirror-dual to fibers with Mordell-Weil torsion of the same degree
- Combinatorial explanation in 2D toric ambient spaces
- We have explicitly checked the conjecture for **all 3145 cases** of codimension two curves
- The combinatorial explanation does not fully carry over to 3D
- We find fibers with new features:
  - Equivalent realizations of the same elliptic curve
  - A fiber with a non-toric zero-section
  - Genus-one fibers with torsion sections in their Jacobian
- We have fully analyzed a self-dual genus one curve that admits quotient and discrete symmetries

## Outlook

- Can we prove the conjecture in general?
- Is there a physical explanation?
- Can we classify discrete symmetries in F-theory via their mirror dual torsion?